

PART VIII

Results based on the NS equations in fully developed turbulence

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- 8 Results based on the equations of the dynamics in fully developed turbulence
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Dynamics based results in fully developed turbulence

Tensorial general expressions

Considering homogeneous isotropic incompressible turbulence.

$$\begin{aligned} Q_i(\mathbf{r}) &= A(r)r_i = \overline{p(\mathbf{x})u_i(\mathbf{x} + \mathbf{r})}_{\text{ex.}} \\ R_{ij}(\mathbf{r}) &= F(r)r_i r_j + G(r)\delta_{ij} = \overline{u_i(\mathbf{x})u_j(\mathbf{x} + \mathbf{r})}_{\text{ex.}} \\ S_{ij\ell}(\mathbf{r}) &= A(r)r_i r_j r_\ell + B(r)(r_i \delta_{j\ell} + r_j \delta_{i\ell}) + D(r)r_\ell \delta_{ij} \\ &= \overline{u_i(\mathbf{x})u_j(\mathbf{x})u_\ell(\mathbf{x} + \mathbf{r})}_{\text{ex.}} \end{aligned}$$

where A, F, G, B, D are arbitrary scalar functions of r^2 ; all even function of r (by isotropy)

- $R_{ij}(\mathbf{r})$ symmetric in i, j : $R_{ij}(\mathbf{r}) = R_{ji}(\mathbf{r})$
- $S_{ij\ell}(\mathbf{r})$ symmetric in i, j : $S_{ij\ell}(\mathbf{r}) = S_{ji\ell}(\mathbf{r})$

General form for HIT

Expression of Q_i

- Continuity condition : $\partial_i u_i = 0$

$$\frac{\partial Q_i(\mathbf{r})}{\partial r_i} = 3A + r \frac{\partial A}{\partial r} = 0, \forall \mathbf{r}$$

$\implies A(r) = 0$ assuming regularity at $r = 0$.

- First order tensor

$$Q_i(\mathbf{r}) \equiv 0, \forall \mathbf{r}.$$

Tensorial general form

Expression of R_{ij}

- Continuity condition :

$$\frac{\partial R_{ij}(\mathbf{r})}{\partial r_i} = r_j \left(4F + r \frac{\partial F}{\partial r} + \frac{1}{r} \frac{\partial G}{\partial r} \right) = 0, \forall \mathbf{r}$$

$$\implies 4F + r \frac{\partial F}{\partial r} + \frac{1}{r} \frac{\partial G}{\partial r} = 0$$

Tensorial general expression

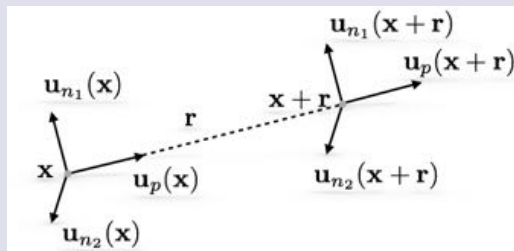
Expression of R_{ij}

- Longitudinal and lateral velocity correlations

$$\overline{u_p(\mathbf{x})u_p(\mathbf{x} + \mathbf{r})} \equiv u^2 f(r)$$

$$\overline{u_n(\mathbf{x})u_n(\mathbf{x} + \mathbf{r})} \equiv u^2 g(r)$$

- $f(r), g(r)$: even scalar functions
- u_{rms} : $u^2 \equiv \overline{u_p^2} = \overline{u_n^2} = \frac{1}{3} \overline{u_i^2}$



General form for HIT

Expression of $S_{ij\ell}$

- From the general form, one obtains

$$s_{ij\ell}(\mathbf{r}) = u^3 \left[\frac{k - rk'}{2r^3} r_i r_j r_\ell + \frac{2k + rk'}{4r} (r_i \delta_{j\ell} + r_j \delta_{i\ell}) - \frac{k}{2r} r_\ell \delta_{ij} \right],$$

where $k(r)$ is the **single scalar function** determining the triple velocity correlation and

$$k'(r) \equiv \frac{\partial k(r)}{\partial r}.$$

- Note that

$$S_{iji}(\mathbf{r}) = \frac{1}{2} u^3 \left[k' + \frac{4}{r} k \right] r_j \equiv \frac{1}{2} K(r) r_j$$

von Kármán equation (1938)

VKH equation

- Write NS equation at points \mathbf{x} and $\mathbf{x}' = \mathbf{x} + \mathbf{r}$.
- Summing and averaging with $u_i \equiv u_i(\mathbf{x}, t)$ and $u'_i \equiv u_i(\mathbf{x} + \mathbf{r}, t)$, in order to write

$$\partial_t \overline{u_i u'_j} = \dots$$

von Kármán equation (1938)

VKH equation

- Previous equation leads to

$$\begin{aligned} \frac{\partial R_{ij}(\mathbf{r})}{\partial t} &= \frac{\partial}{\partial r} \left[\overline{u_i u_k u'_j} - \overline{u_i u'_k u_j} \right] \\ &+ \left[\frac{\partial}{\partial r_i} \overline{p u'_j} - \frac{\partial}{\partial r_j} \overline{p' u_i} \right] \\ &+ 2\nu \frac{\partial^2 R_{ij}(\mathbf{r})}{\partial r_\ell \partial r_\ell} \end{aligned}$$

von Kármán equation (1938)

VKH equation

- Putting $i = j$

$$\frac{\partial R(r)}{\partial t} = \frac{1}{2} \left(r \frac{\partial}{\partial r} + 3 \right) K(r) + 2\nu \left(\frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} \right) R(r)$$

- First integral of this equation gives

$$\frac{\partial(u^2 f)}{\partial t} = \left(\frac{\partial}{\partial r} + \frac{4}{r} \right) (u^3 k) + 2\nu \left(\frac{\partial^2}{\partial r^2} + \frac{4}{r} \frac{\partial}{\partial r} \right) (u^2 f)$$

Kolmogorov 4/5 law

Structure functions

Considering the velocity structure functions

- 2nd order

$$B_{ik}(\mathbf{r}) = \overline{\delta u_i(\mathbf{r}) \delta u_k(\mathbf{r})} = \overline{(u_i(\mathbf{x} + \mathbf{r}) - u_i(\mathbf{x})) (u_k(\mathbf{x} + \mathbf{r}) - u_k(\mathbf{x}))}$$

- 3rd order

$$\begin{aligned} B_{ik\ell}(\mathbf{r}) &= \overline{\delta u_i(\mathbf{r}) \delta u_k(\mathbf{r}) \delta u_\ell(\mathbf{r})} \\ &= \overline{(u_i(\mathbf{x} + \mathbf{r}) - u_i(\mathbf{x})) (u_k(\mathbf{x} + \mathbf{r}) - u_k(\mathbf{x})) (u_\ell(\mathbf{x} + \mathbf{r}) - u_\ell(\mathbf{x}))} \end{aligned}$$

And replacing in the von Kármán Howarth equation yields

$$-\frac{2}{3}\varepsilon + \frac{1}{2} \frac{\partial B_{pp}}{\partial t} = \frac{1}{6r^4} \frac{\partial}{\partial r} (r^4 B_{ppp}) - \frac{\nu}{r^4} \frac{\partial}{\partial r} \left(r^4 \frac{\partial}{\partial r} B_{pp} \right)$$

with

$$\frac{\partial E}{\partial t} = \frac{1}{2} \frac{\partial}{\partial t} \overline{u_i^2} = -\varepsilon = \frac{3}{2} \frac{\partial}{\partial t} \overline{u_p^2}.$$

Kolmogorov 4/5 law

VKH to 4/5 law

- Neglecting the time derivative term compared to ε .
- Neglecting the dissipative term in the inertial range
- Integrating over r
- One obtains in the inertial Range

$$B_{ppp} = \overline{(\delta u_p(\mathbf{r}))^3} = -\frac{4}{5}\varepsilon r$$

- But observed only for very large Reynolds number!!!

Bibliography

Books

- Batchelor 1953, The Theory of Homogeneous Turbulence
- Frisch 1995, Turbulence
- Pope 2000, Turbulent Flow
- Monin & Yaglom, Statistical Fluids Mechanics, Mechanics of Turbulence

Final Exam

Objective

- Start by determining the relation, equivalent to the '4/5 law', for the burgers equations which is written for $\langle \delta u^3 \rangle$. A rigorous demonstration is expected.
- Extend this demonstration to the case of 3D turbulent flow by writing the von Karman equation and to the famous Kolmogorov 4/5 law.
- Your mission consists in giving all the details and assumptions of the derivation.
- Write the demonstration for the equivalent law in the case of the passive scalar dynamics.